Mass of homogeneous sphere: 
$$M = \rho V = \frac{4}{2}\pi\rho r^3$$
 [1]

Schwarzschild radius: 
$$r_{S} = \frac{2GM}{c^{2}} \div \frac{r_{S}}{r} = \frac{2GM}{rc^{2}} = \frac{2G}{rc^{2}} \div \frac{4}{3}\pi\rho r^{3} = \frac{8\pi G\rho}{3c^{2}}r^{2}$$
[2]

$$v_{esc} = \sqrt{\frac{2GM}{r}} \therefore v_{esc}^2 = \frac{2GM}{r} = \frac{r_s c^2}{r} \therefore \frac{r_s}{r} = \frac{v_{esc}^2}{c^2}$$
[3]

$$\gamma = \frac{1}{\sqrt{1 - \frac{\nu^2}{c^2}}} \tag{4}$$

Relativistic mass:
$$m = \gamma m_0$$
[5]Kinetic energy: $E_k = mc^2 - m_0c^2 = \gamma m_0c^2 - m_0c^2 = (\gamma - 1)m_0c^2$ [6]

$$\therefore \ \frac{E_k}{m_0 c^2} = \gamma - 1 \therefore 1 + \frac{E_k}{m_0 c^2} = \gamma \therefore \gamma = \frac{m_0 c^2 + E_k}{m_0 c^2}$$
[7]

from [2]:

Escape velocity:

Lorentz factor:

$$\therefore \quad \sqrt{1 - \frac{v^2}{c^2}} = \frac{m_0 c^2}{m_0 c^2 + E_E}$$
[8]

radius: 
$$\sqrt{1 - \frac{v_{esc}^2}{c^2}} = \frac{m_0 c^2}{m_0 c^2 + E_{esc}}$$
 [9]

using [3]: 
$$\therefore \quad \sqrt{1 - \frac{r_s}{r_{esc}}} = \frac{m_0 c^2}{m_0 c^2 + E_{esc}} \therefore 1 - \frac{r_s}{r_{esc}} = \left(\frac{m_0 c^2}{m_0 c^2 + E_{esc}}\right)^2$$
[10]

$$\therefore \frac{r_s}{r_{scc}} = 1 - \left(\frac{m_0 c^2}{m_0 c^2 + E_{scc}}\right)^2$$
[11]

$$\frac{r_{esc}}{r} = \frac{8\pi G\rho}{3c^2} r_{esc}^2$$
[12]

$$\therefore \ r_{esc}^2 = \frac{3c^2}{8\pi G\rho} \cdot \left[ 1 - \left( \frac{m_0 c^2}{m_0 c^2 + E_{esc}} \right)^2 \right]$$
[13]

$$\therefore r_{esc} = c \cdot \sqrt{\frac{3}{8\pi G\rho}} \cdot \sqrt{1 - \left(\frac{m_0 c^2}{m_0 c^2 + E_{esc}}\right)^2}$$
[14]

An object with rest mass  $m_0$  and kinetic energy  $E_{esc}$  can/will escape if it is at a distance of at least  $r_{esc}$ .

For a neutron star, we've got: 
$$\rho = \rho_N = 1.392134 \times 10^{18} \text{ kg/m}^3$$
 [15]

and of course: 
$$c = 299792458 \text{ m/s}$$
 [16]

as well as:

so:

$$G = 6.67408 \times 10^{-11} \text{ m}^3/\text{kg/s}^3$$
 [17]

so:  $c \cdot \sqrt{\frac{3}{8\pi G\rho}} = 299792458 \cdot \sqrt{\frac{3}{8\pi \cdot 6.67408 \times 10^{-11} \cdot 1.392134 \times 10^{18}}} = 10745 \text{ m}$  [18]

This value equals the critical black hole radius as calculated in <u>http://henk-reints.nl/astro/HR-on-the-universe.php</u>.

Suppose this object is an electron originating from a decaying neutron of a neutron star without obtaining *energy* from the star's *binding energy*. Then:

Electron mass: 
$$m_0 = m_e = 9.10938356 \times 10^{-31} \text{ kg}$$
 [19]

$$m_0 c^2 = m_e c^2 = 8.18710565 \times 10^{-14} \,\text{J}$$
 [20]

Neutron decay *energy*: 
$$E_{dec} = (m_n - m_p - m_e)c^2 = 0.78257 \text{ MeV} = 1.2538 \times 10^{-13} \text{ J}$$
 [21]

First assume the electron obtains all this *energy*. Then:

$$r_0 = 10745 \cdot \sqrt{1 - \left(\frac{8.18710565 \times 10^{-14}}{8.18710565 \times 10^{-14} + 1.2538 \times 10^{-13}}\right)^2} = 9871 \,\mathrm{m}$$
 [22]

$$M_0 = \frac{4}{3}\pi\rho r_0^3 = 7.234 \times 10^{30} \,\mathrm{kg} \approx 3.637 M_{\odot}$$
[23]

So a neutron star with a diameter that exceeds 19.74 km and thus a *mass* above 3.637 times that of our sun cannot lose any electrons by neutron decay because *gravitation* definitely wins that game. Above this size a ball of neutronium would not be able to disintegrate. That's why I do still not really understand what the big bang actually was. The *IniAll* must have had a *radius* of nearly 3 astronomical units.

In practice, the electron antineutrino takes away roughly 50% of the neutron decay energy on average,

$$r_1 = 10745 \cdot \left( \frac{8.18710565 \times 10^{-14}}{8.18710565 \times 10^{-14} + \frac{1.2538 \times 10^{-13}}{2}} \right)^2 = 8856 \text{ m}$$
 [24]

corresponding to a *diameter* of: and its corresponding *mass* is:

$$d_1 = 17.71 \text{ km}$$
[25]  
$$M_1 = \frac{4}{3}\pi\rho r_1^3 = 4.050 \times 10^{30} \text{ kg} \approx 2.036 M_{\odot}$$
[26]

To me it is remarkable that this *mass* is within the tolerance of the lower bound of the TOV-limit as stated on <u>https://en.wikipedia.org/wiki/Tolman%E2%80%93Oppenheimer%E2%80%93Volkoff\_limit#History</u> (as of 2019-08-27): The mass of the pulsar PSR J0348+0432, at 2.01±0.04 solar masses, puts an empirical lower bound on the TOV limit.

Below this *mass*, a neutron star would be prone to evaporation by neutron decay. This could mean twice the solar *mass* is a lower *mass* limit for neutron stars to be able to stay alive at all. But the neutrons in a neutron star are definitely not free neutrons, so they may have a far longer *decay time*, which would also avoid this evaporation.